

# CONFORMAL HARMONIC ANALYSIS AND INTERTWINING DIFFERENTIAL OPERATORS

H.P. Jakobsen

## INTRODUCTION

Earlier in this conference, we have heard Professor I.E. Segal's lecture "Causality and Symmetry in Cosmology and the Conformal Group". I refer to that for motivation, general background, and references. Here shall be described a few results concerning the development of the theory towards elementary particle physics; in particular some results of a study, jointly with Michele Vergne (1), of the interplay between representations of the conformal group (in the form of  $SU(2,2)$ ) and classical mass-zero-equations, will be given.

## GENERALITIES

For technical reasons it is convenient to work with the group

$$G = \{g \in SL(4, \mathbb{C}) \mid g^* \begin{pmatrix} 0 & \mathbf{i} \\ -\mathbf{i} & 0 \end{pmatrix} g = \begin{pmatrix} 0 & \mathbf{i} \\ -\mathbf{i} & 0 \end{pmatrix}\}.$$

This group is isomorphic to the group

$$SU(2,2) = \{g \in SL(4, \mathbb{C}) \mid g^* \begin{pmatrix} \mathbf{I} & 0 \\ 0 & -\mathbf{I} \end{pmatrix} g = \begin{pmatrix} \mathbf{I} & 0 \\ 0 & -\mathbf{I} \end{pmatrix}\}$$

which is used in Segal's theory, and it is straightforward to translate results from  $G$  to  $SU(2,2)$ .

$G$  acts on the generalized upper half-plane  $\mathcal{D} = \{2 \times 2 \text{ complex matrices } z \mid \frac{z-z^*}{2i} \text{ is positive definite}\}$  by, for  $G \ni g = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ ;  $a, b, c, d, 2 \times 2$  complex matrices,

$$g \cdot z = \frac{az+b}{cz+d}.$$

One can prove that  $(g \cdot z^*)^* = g \cdot z$ , and thus  $G$  leaves the boundary of  $\mathcal{D}$ ,  $\partial\mathcal{D} = H(2) = \{2 \times 2 \text{ complex matrices } h \mid h = h^*\}$ , invariant, and this gives rise to the well-known local action of  $SU(2,2)$  on Minkowski space  $M$ , identified with  $H(2)$  by

$$M \ni (t, x, y, z) \iff h = \begin{pmatrix} t+z & x+iy \\ x-iy & t-z \end{pmatrix} \in H(2).$$

Finally, let  $c: \begin{pmatrix} \lambda_1 & \lambda_2 \\ \lambda_3 & \lambda_4 \end{pmatrix} \rightarrow \begin{pmatrix} \lambda_4 & -\lambda_2 \\ -\lambda_3 & \lambda_1 \end{pmatrix}$  denote the co-factor mapping on the space of  $2 \times 2$  complex matrices and observe that  $c$ 's restriction to  $H(2)$  is the operation of space-reversal.

### THE REPRESENTATIONS

Let  $\psi$  be a function from  $H(2)$  to  $\mathbb{C}$ ,  $g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in G$ , and  $n \in \mathbb{Z}$  and define

$$(U_n(g)\psi)(h) = \det(ch+d)^{-(n+2)} \psi(g^{-1}h).$$

Moreover, let  $\varphi$  be a function from  $H(2)$  to  $\mathbb{C}^2$  and define

$$(D_n^+(g)\varphi)(h) = (ch+d)^{-1} \det(ch+d)^{-(n+2)} \varphi(g^{-1}h),$$

$$(D_n^-(g)\varphi)(h) = (hc^*+d^*) \det(hc^*+d^*)^{-(n+3)} \varphi(g^{-1}h),$$

and finally define, on functions from  $H(2)$  to  $\mathbb{C}^4$ ,

$$D_n(g) = D_n^+(g) \oplus D_n^-(g).$$

$U_n$  and  $D_n$  are then well-defined actions on measurable functions.

It is proved in (1), but has also been observed elsewhere (see References in (1)), that one has

*Theorem 1.* For  $n \geq 0$  there exist Hilbert spaces  $H_n$  and  $K_n$ , obtained as completions of spaces of positive energy and positive mass functions, such that  $U_n$  is unitary and irreducible in  $H_n$ , and  $D_n$  is the direct sum of two unitary irreducible representations in  $K_n$ . Moreover, there exists a Hilbert space  $H_{-1}$  ( $K_{-1}$ ) consisting of solutions to  $\square\psi = 0$  ( $\not\psi = 0$ ) in which  $U_{-1}$  ( $D_{-1}$ ) is unitary.

The significance of the representations  $U_n$  and  $D_n$  for integers  $n \leq -1$  and the relation of these to positive mass functions is illustrated by

THE INTERTWINING RELATIONS

For the sake of being explicit, let

$$\square = \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2},$$

$$\sigma = \begin{pmatrix} \frac{\partial}{\partial t} + \frac{\partial}{\partial z} & \frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} - i \frac{\partial}{\partial y} & \frac{\partial}{\partial t} - \frac{\partial}{\partial z} \end{pmatrix}, \quad \tilde{\sigma} = \begin{pmatrix} \frac{\partial}{\partial t} - \frac{\partial}{\partial z} & -\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \\ -\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} & \frac{\partial}{\partial t} + \frac{\partial}{\partial z} \end{pmatrix},$$

and  $\not\psi = \begin{pmatrix} 0 & \sigma \\ \tilde{\sigma} & 0 \end{pmatrix}$ . (Then  $\not\psi^2 = \square$ .)

*Theorem 2.* (1)  $\forall g \in G$ :

$$\square^n U_{-n}(g) = U_n(g) \square^n$$

$$\not\psi^{2n+1} D_{-1-n}(g) = D_n(g) \not\psi^{2n+1}$$

whenever both sides make sense.

Let, for  $n \in \mathbb{N}$ , two functions  $\psi_i: H(2) \rightarrow \mathbb{C}$  ( $i=1,2$ ) be

equivalent if  $\square^n(\psi_1 - \psi_2) = 0$  and denote the equivalence classes by  $[\cdot]_n^0$ . Likewise define, for non-negative integers  $n$ , two functions  $\varphi_i: H(2) \rightarrow \mathbb{C}^4$  ( $i=1,2$ ) to be equivalent if  $\nabla^{2n+1}(\varphi_1 - \varphi_2) = 0$  and denote the equivalence classes by  $[\cdot]_n^{\frac{1}{2}}$ .

*Theorem 3.* (1) It is possible to put a Hilbert space structure  $H_n(K_n)$  on a space of equivalence classes  $[\cdot]_n^0([\cdot]_n^{\frac{1}{2}})$  of positive mass function, such that  $U_{-n}(D_{-1-n})$  is a unitary representation in  $H_n(K_n)$ .

*Remark.* If one restricts the representations  $U_n$  and  $D_n$  to the Poincaré group together with the scale-transformations, the proofs of the above mentioned facts are quite simple. Also, theorem 3 can then be sharpened, in the sense that for each  $U_{-n}(D_{-1-n})$  ( $n$  positive (or zero)) there is a natural Hilbert space  $H_{-n}(K_{-n})$  of positive mass functions, on which this restriction is unitary. However, this is not the case for the full group  $G$ . In fact, for the one-parameter group

$$\tau(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \subseteq G$$

(which is the  $G$ -version of the uni-time translation subgroup of  $SU(2,2)$ ),  $U_{-n}(\tau(\theta))(D_{-1-n}(\tau(\theta)))$  does not leave  $H_{-n}(K_{-n})$  invariant, as can be seen easily by computations on special functions.

#### REFERENCE

1. H.P. JAKOBSEN and M. VERGNE, Wave and Dirac operators, and representations of the conformal group. (To appear in J. Functional Analysis.)