

GROUP THEORETICAL ASPECTS OF THE CHRONOMETRIC THEORY.*

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Below is given a brief outline of one particular point, namely that of chronometric quantum-numbers for mass-zero equations. A complete description of these and related matters, which results from the joint investigations of I.E. Segal, B. Speh, B. Ørsted and myself, will appear elsewhere. In this connection I would like to thank Segal, Speh, Ørsted and M. Harris for valuable discussions. The other aspects touched upon in the talk are described in [1], and [2]. The chronometric theory is described in [5]. For additional background and further results, c.f. [3] and literature cited there.

1. Notation

$\tau = \tau_1$ denotes the defining representation of $GL(2, \mathbb{C})$, i.e. $\tau \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}$. τ_n denotes the n-th fold symmetrized tensor product of τ , acting on $V_n = \bigotimes_s^n \mathbb{C}^2$, and τ_0 denotes the trivial representation. The letters a, b, c, d, h, k, u, v, z, w denote 2×2 matrices.

2. Basics

The unitary representations of $SU(2,2)$ corresponding to spin $\frac{n}{2}$ and mass 0 are

$$\begin{aligned} (U_n^+(g)f)(z) &= \tau_n(cz + d)^{-1} \det(cz + d)^{-1} f(g^{-1}z), \text{ and} \\ (U_n^-(g)f)(z) &= \tau_n(a - (g^{-1}z)c)^{-1} \det(a - (g^{-1}z)c)^{n+1} f(g^{-1}z), \end{aligned} \tag{2.1}$$

where $g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SU(2,2)$, $g^{-1}z = \frac{az+b}{cz+d}$, and $f(z) \in V_n$.

* This research was supported in part by a grant from the N.S.F.

Let

$$SU(2,2)_D = \left\{ g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid g^* \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} g = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \right\},$$

$$SU(2,2)_B = \left\{ g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid g^* \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} g = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right\},$$

$$D = \left\{ z \mid \frac{z-z^*}{2i} \text{ is strictly positive definite} \right\}, \text{ and}$$

$$B = \left\{ z \mid z^* z < 1 \right\}.$$

Of course, the groups $SU(2,2)_D$ and $SU(2,2)_B$ are isomorphic. D is the generalized upper half-plane and B is the generalized unit disk. The Shilov boundary of D is $H(2) = \{ 2 \times 2 \text{ hermitian matrices} \} = \text{Minkowski space}$, whereas the Shilov boundary of B is $U(2) = \{ 2 \times 2 \text{ unitary matrices} \}$. $\widetilde{U(2)}$ = the universal covering group of $U(2) = \mathbb{R} \times SU(2) = \mathbb{R} \times S^3$ is the Segal cosmos.

The formulas (2.1) define representations of $SU(2,2)_D$ as well as $SU(2,2)_B$ provided that for $SU(2,2)_D$, f is a holomorphic function on D , and that for $SU(2,2)_B$, f is a holomorphic function on B . These representations are equivalent. In both versions one can pass to the Shilov boundary and thus realize the representations on the respective space-times. One reason that the following claims are true is that the transition from $SU(2,2)_D$ to $SU(2,2)_B$ (and from D to B) can be obtained through the action of an element in the complexification of $SU(2,2)$, and that all the expressions are analytic in the variables g and z . Observe that once a choice of domain has been made the term $(a - (g^{-1}z)c)^{-1}$ may be simplified: For $g^{-1} \in SU(2,2)_D$ and $z \in D$ it is equal to $zc^* + d^*$ and in the B -situation it is equal to $(a^* + zb^*)$.

3. Chronometric quantum numbers

The problem may be handled by means of the "ladder representations" [4] (see also [1]). In contrast, the approach presented here is intrinsic in the sense that it deals directly with the relevant function spaces and differential equations. In the following we shall only consider the representations U_n^+ since the treatment of the U_n^- 's follows from this by obvious modifications.

The domain D is particularly nice for describing the Hilbert space H_n^+ that carries the unitary irreducible representation U_n^+ : H_n^+ is a reproducing kernel Hilbert space and in the D -version the kernel is given as

$$K_n^+(z, w) = \zeta_n \left(\frac{z-w^*}{2i} \right)^{-1} \det \left(\frac{z-w^*}{2i} \right)^{-1} = \int_{b(C^+)} \zeta_n(k) e^{i \operatorname{tr}(z-w^*)k} dm(k), \quad (3.1)$$

where $b(C^+) = \{k = k^* \mid \operatorname{tr} k \geq 0 \text{ and } \det k = 0\}$ is the boundary of the solid forward light-cone C^+ , and $dm(k)$ is the usual Lorentz-invariant measure.

We wish to determine a complete set of quantum numbers for the representations U_n^+ . From the point of view of Segal's cosmological theory as well as from that of group theory, the natural way to do this is to use the maximal compact subgroup K of $SU(2,2)$. K coincides with the maximal subgroup of isometries of $SU(2,2)$ under its action on $U(2)$. In terms of group theory we shall thus determine the decomposition of the restriction of U_n^+ to K as a direct sum of finite-dimensional representations of K ; the so-called K -types.

In $SU(2,2)_B$ we thus choose K as

$$K = \left\{ \begin{pmatrix} u & 0 \\ 0 & v \end{pmatrix} \mid u, v \in U(2), \det(u \ v) = 1 \right\} \quad (3.2)$$

In particular, $K = U(1) \times SU(2) \times SU(2)$.

$$T = \left\{ \begin{pmatrix} e^{i\tau/2} & 0 \\ 0 & e^{-i\tau/2} \end{pmatrix} \mid \tau \in \mathbb{R} \right\} \text{ is Segal's time-translation subgroup of } K.$$

The natural domain to use for the determination of the K -types is B . On B ,

$$(U_n^+(u, v)f)(z) = (U_n^+ \begin{pmatrix} u & 0 \\ 0 & v \end{pmatrix} f)(z) = \zeta_n(v) \det v f(u^{-1}z \ v), \quad (3.3)$$

hence the finite-dimensional K -irreducible subspaces consist of V_n -valued polynomials.

(3.3) clearly defines an action of K on the set of all V_n -valued polynomials on B .

We denote this action by $\widetilde{U}_n^+$.

Consider the representation P of K on the space of all $\mathbb{C}$ -valued polynomials defined by

$$(P(u, v)p)(z) = p(u^{-1}z \ v) \quad (3.4)$$

$$\text{Proposition 3.1} \quad P(u, v) = \bigoplus_{l=0}^{\infty} \bigoplus_{r=0}^{\infty} \frac{\zeta_1(v) \otimes \zeta_1(\bar{u})}{\det u^{2r}}$$

For later use we observe that P leaves the ideal $I = I(\det z)$ of polynomials propor-

tional to $\det z$, invariant. Clearly then

$$\text{Proposition 3.2} \quad P(u, v) \Big|_I = \bigoplus_{l=0}^{\infty} \bigoplus_{r=1}^{\infty} \frac{\tau_l(v) \otimes \tau_l(\bar{u})}{\det u^{2r}}$$

The K-types of U_n^+ are, of course, to be found among the summands of $\widetilde{U}_n^+(u, v) = \tau_n(v) \det v \otimes P(u, v)$.

Proposition 3.3

$$\tau_n(v) \det v \otimes P(u, v) = \bigoplus_{l=0}^{\infty} \bigoplus_{r=0}^{\infty} \bigoplus_{\gamma=0}^{\min\{n, l\}} \frac{\tau_{n+l-2\gamma}(v) \otimes \tau_l(\bar{u})}{\det u^{\gamma+1+2r}}$$

However, since the Hilbert space of U_n^+ consists of solutions to certain differential equations, some of the summands must be excluded. To see exactly which do not occur we return to the domain D . By expanding the functions on D around the point $i = \begin{pmatrix} i & 0 \\ 0 & i \end{pmatrix}$ and by choosing K in this version as the subgroup of $SU(2, 2)_D$ that leaves i fixed, it follows (e.g. by a Gram-Schmidt argument) that the finite-dimensional subspaces of functions on D that carry the irreducible K-representations are built up from functions of the form

$$F_{\underline{q}}(z) = \int_{b(C^+)} \tau_n(k) e^{i \operatorname{tr}(z+i)k} \underline{q}(k) \, dm(k), \quad (3.5)$$

where $\underline{q}$ is a V_n -valued polynomial. However, $\underline{q}$ determines a non-zero function $F_{\underline{q}}$ if and only if there is an x in V_n such that $\langle \tau_n(k) \underline{q}(k), x \rangle$ is not identically zero on $b(C^+)$. That is: Exactly those $\underline{q}$'s for which for all x in V_n , $\langle \tau_n(k) \underline{q}(k), x \rangle = 0$ on $b(C^+)$ do not contribute to the K-types in U_n^+ . Now observe that $b(C^+)$ is a sufficiently big subset of the variety of the prime ideal $I(\det z)$ that we may conclude:

$$\forall x \in V_n : \langle \tau_n(k) \underline{q}(k), x \rangle = 0 \text{ on } b(C^+) \iff$$

$$\forall x \in V_n : \langle \tau_n(z) \underline{q}(z), x \rangle \in I(\det z).$$

$$\text{Let } A = \{ \underline{q} \mid \forall x \in V_n : \langle \tau_n(z) \underline{q}(z), x \rangle \in I(\det z) \}$$

A is clearly invariant under the action of $\widetilde{U}_n^+$, and even though it originally was determined from the domain D , it follows by analyticity that

Proposition 3.4 The K-types in A are exactly those in $\widetilde{U}_n^+$ that do not occur in U_n^+ .

Let $\mathcal{S}$ denote the linear map of $A \otimes V_n$ into the space of $\mathbb{C}$ -valued polynomials,

$$(\mathcal{S}(\underline{q} \otimes x))(z) = \langle \tau_n(z) \underline{q}(z), \bar{x} \rangle \quad (3.7)$$

Then

$$P(u, v) \mathcal{S} = \mathcal{S}(\tilde{U}_n^+(u, v) \otimes \det \bar{v} \tau_n(\bar{u})) \quad (3.8)$$

Now choose a summand $\frac{\tau_{n+1-2r}(v) \otimes \tau_1(\bar{u})}{\det u^{r+1+2r}}$

from $\tilde{U}_n^+$ and assume that $\underline{q} \in A$ transforms according to this representation. It follows from (3.8) and Proposition 3.1 that $\langle \tau_n(z) \underline{q}(z), \bar{x} \rangle$ transforms according to

$$\frac{\tau_{n+1-2r}(v) \otimes \tau_{n+1-2r}(\bar{u})}{\det u^{2r+2r}}.$$

Since $\underline{q} \in A$ it follows from Proposition 3.2 that either $r > 0$ or $r = 0$. Thus

Proposition 3.5

$$U_n^+(u, v) = \bigoplus_{l=0}^{\infty} \frac{\tau_{n+1}(v) \otimes \tau_l(\bar{u})}{\det u}$$

By analogous reasoning it follows that

Proposition 3.6

$$U_n^-(u, v) = \bigoplus_{l=0}^{\infty} \frac{\tau_l(v) \otimes \tau_{l+n}(\bar{u})}{\det u}$$

REFERENCES

- [1] Jakobsen, H.P. and Vergne, M., J. Functional Analysis 24, 52 (1977); 34, 29 (1979).
- [2] Jakobsen, H.P., in "Non-Commutative Harmonic Analysis", Lecture Notes in Math. 728, Berlin-Heidelberg-New York: Springer Verlag 1979.
- [3] Jakobsen, H.P., Ørsted, B., Segal, I.E., Speh B., and Vergne, M., Proc. Nat. Acad. Sci. USA 75, 1609 (1978).
- [4] Mack, G. and Todorov, I., J. Math. Phys. 10, 2078 (1969).
- [5] Segal, I.E., "Mathematical Cosmology and Extragalactic Astronomy", Academic Press, New York, 1976.