

The equivalence of two parametrizations for the $I(2)$ model and an example of the conditions for asymptotic χ^2 inference

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1 The parametrization of Boswijk.

We here discuss the parametrization introduced by Boswijk, which normalizes β on a matrix c , and η on a matrix d . We first prove the following lemma

Lemma 1 *Without loss of generality we can assume that the parameters satisfy the normalizations*

$$c'\beta = I_r, c'_\perp\beta_\perp = I_{p-r}, d'\eta = I_s, d'_\perp\eta_\perp = I_{p-r-s}, \quad (1)$$

where the normalization matrices satisfy the conditions

$$c'c = I_r, c'_\perp c_\perp = I_{p-r}, d'd = I_s, d'_\perp d_\perp = I_{p-r-s}. \quad (2)$$

PROOF. We choose a value of the parameters $\alpha^0, \beta^0, \Gamma^0, \xi^0, \eta^0$ satisfying the $I(2)$ condition. Because $\text{rank}(\beta^0) = r$, we can find a permutation matrices S , so that with $c' = (I_r, 0)S$ we have $|c'\beta^0| \neq 0$. Similarly we can find a permutation matrix S_0 so that with $d' = (I_s, 0)S_0$ we get $|d'\eta^0| \neq 0$. For this choice of S and S_0 , consider the subset

$$\Theta_0 = \{\beta, \eta, \dots : |c'\beta| \neq 0, |d'\eta| \neq 0\}.$$

This is not the full parameter space, but the complement of Θ_0 , where either $|c'\beta| = 0$ or $|d'\eta| = 0$, has Lebesgue measure zero. Thus "without loss of generality" we can take $(\beta, \eta) \in \Theta_0$.

We further define

$$c_\perp = S' \begin{pmatrix} 0 \\ I_{p-r} \end{pmatrix}, d_\perp = S'_0 \begin{pmatrix} 0 \\ I_{p-r-s} \end{pmatrix}$$

so that (2) holds, and apply these to define the normalized parameters

$$\beta_c = \beta(c'\beta)^{-1}, \alpha_c = \alpha\beta'c, \eta_d = \eta(d'\eta)^{-1}, \xi_d = \xi\eta'd, \\ \beta_{c_\perp} = (I_p - c\beta'_c)c_\perp, \eta_{d_\perp} = (I_{p-r} - d\eta'_d)d_\perp$$

so that (1) holds. This completes the proof of Lemma 2 ■

In the following we assume that Lemma 1 is satisfied and call the parameters $\alpha_c, \beta_c, \Gamma, \xi_d$, and η_d , to indicate the way they are normalized. We let $c_1 = c_\perp d$ and $c_2 = c_\perp d_\perp$ and introduce the parameters A_i , see Boswijk (2000, page 898), as

$$A_1 = -\beta'_c c_1, \quad A_2 = -\beta'_c c_2, \quad A_3 = -\bar{\alpha}'_c \Gamma \beta_{c2}, \quad A_4 = -\eta'_d d_\perp.$$

The relation between the parameters $\beta_c, \beta_{c\perp}, \eta_d, \eta_{d\perp}$, and the parameters A_i is given by

$$\begin{aligned} \beta'_c &= c' - A_1 c'_1 - A_2 c'_2 = c' - (A_1 d' + A_2 d'_\perp) c'_\perp, \\ \beta_{c\perp} &= c_\perp + c(A_1 d' + A_2 d'_\perp), \\ \eta'_d &= d' - A_4 d'_\perp, \\ \eta_{d\perp} &= d_\perp + d A_4, \end{aligned} \tag{3}$$

from which we find

$$\begin{aligned} \beta_{c2} &= \beta_{c\perp} \eta_{d\perp} = c(A_1 A_4 + A_2) + c_1 A_4 + c_2, \\ \beta_{c1} &= \beta_{c\perp} (\beta'_{c\perp} \beta_{c\perp})^{-1} \eta_d = (c_\perp + c(A_1 d' + A_2 d'_\perp)) (\beta'_{c\perp} \beta_{c\perp})^{-1} \eta_d, \\ \tau_c &= \beta_{2c\perp} = (c - c_2(A'_4 A'_1 + A'_2), c_1 - c_2 A'_4). \end{aligned} \tag{4}$$

In the present paper we have normalized the parameters differently. We assume that β and τ are normalized on $\bar{\beta}^0$ and $\bar{\tau}^0$. Thus our parameters can be recovered from α_c, β_c, \dots or A_i as follows:

$$\begin{aligned} \beta &= \beta_c (\bar{\beta}_c^{0'} \beta_c)^{-1} = (c - c_1 A'_1 - c_2 A'_2) (\bar{\beta}_c^{0'} \beta_c)^{-1}, \\ \tau &= \tau_c (\bar{\tau}_c^{0'} \tau_c)^{-1} = (c - c_2(A'_4 A'_1 + A'_2), c_1 - c_2 A'_4) (\bar{\tau}_c^{0'} \tau_c)^{-1}. \end{aligned}$$

Note that $\beta^0 = \beta_c^0$, and we choose $\beta_1 = \beta_{c1}$ and $\beta_2 = \beta_{c2}$. We can now prove

Lemma 2 *The first order approximation of the parameters (B_0, B_1, B_2, C) in terms of A_i are given by*

$$\begin{aligned} B_0 &\sim (\beta_2^{0'} \beta_2^0)^{-1} [(A'_3 - A_3^{0'}) - (A'_4 - A_4^{0'}) (A_1^{0'} c' + c'_1) \psi^0] \\ B_1 &\sim -A'_1 + A_1^{0'} \\ B_2 &\sim -(\beta_2^{0'} \beta_2^0)^{-1} (A_4^{0'} (A'_1 - A_1^{0'}) + (A'_2 - A_2^{0'})) \\ C &\sim -(\beta_2^{0'} \beta_2^0)^{-1} (A'_4 - A_4^{0'}) \rho_\perp^{0'} \rho_\perp^0 \end{aligned} \tag{5}$$

where \sim indicates that terms with higher order have been dropped.

PROOF. We find from (36)

$$\begin{aligned} B_2 &= \bar{\beta}_2^{0'} (\beta - \beta^0) = -(\beta_2^{0'} \beta_2^0)^{-1} (\beta_2^{0'} c_1 (A'_1 - A_1^{0'}) + \beta_2^{0'} c_2 (A'_2 - A_2^{0'})) (\bar{\beta}_c^{0'} \beta_c)^{-1} \\ &= -(\beta_2^{0'} \beta_2^0)^{-1} (A_4^{0'} (A'_1 - A_1^{0'}) + (A'_2 - A_2^{0'})) (\bar{\beta}_c^{0'} \beta_c)^{-1}, \end{aligned}$$

using the relations $\beta_2^{0'} c_2 = I_{p-r-s}$ and $\beta_2^{0'} c_1 = A_4^{0'}$. It is seen that the parameter $(A_4^{0'} (A'_1 - A_1^{0'}) + (A'_2 - A_2^{0'})) (\bar{\beta}_c^{0'} \beta_c)^{-1}$ is the non-linear combination of A_i parameters

which is equivalent to B_2 , which is normalized by T^2 when estimated. Because $\bar{\beta}_c^{0'}\beta_c \sim I_r$ close to β^0 we find

$$B_2 \sim -(\beta_2^{0'}\beta_2^0)^{-1}(A_4^{0'}(A'_1 - A_1^{0'}) + (A'_2 - A_2^{0'}))$$

below we apply this result to replace $A'_2 - A_2^{0'}$ by $-A_4^{0'}(A'_1 - A_1^{0'})$ in the expansions.

Similarly we find

$$B_1 = \bar{\beta}_1^{0'}(\beta - \beta^0) = -(\beta_1^{0'}\beta_1^0)^{-1}(\beta_1^{0'}c_1(A'_1 - A_1^{0'}) + \beta_1^{0'}c_2(A'_2 - A_2^{0'}))(\bar{\beta}_c^{0'}\beta_c)^{-1}.$$

replacing $A'_2 - A_2^{0'}$ by $-A_4^{0'}(A'_1 - A_1^{0'})$ and $\bar{\beta}_c^{0'}\beta_c$ by I_r , we find

$$B_1 \sim -(\beta_1^{0'}\beta_1^0)^{-1}(\beta_1^{0'}c_1 - \beta_1^{0'}c_2\beta_2^{0'}c_1)(A'_1 - A_1^{0'})$$

From (4) we find

$$\beta_1^{0'}c_1 = \eta^{0'}(\beta_{\perp}^{0'}\beta_{\perp}^0)^{-1}d, \quad \beta_1^{0'}c_2 = \eta^{0'}(\beta_{\perp}^{0'}\beta_{\perp}^0)^{-1}d_{\perp}, \quad \beta_2^{0'}c_1 = A_4^{0'},$$

so that the coefficient matrix becomes

$$\eta^{0'}(\beta_{\perp}^{0'}\beta_{\perp}^0)^{-1}(d - d_{\perp}A_4^{0'}) = \eta^{0'}(\beta_{\perp}^{0'}\beta_{\perp}^0)^{-1}\eta^0 = \beta_1^{0'}\beta_1^0,$$

so that with $\bar{\beta}_c^{0'}\beta_c \sim I_r$ we find

$$B_1 \sim -A'_1 + A_1^{0'}.$$

Next we find an expression for C :

$$\begin{aligned} C &= \bar{\beta}_2^{0'}(\tau - \tau^0)\rho_{\perp} = \bar{\beta}_2^{0'}\tau_c(\bar{\tau}^{0'}\tau_c)^{-1}\rho_{\perp} \\ &= -(\beta_2^{0'}\beta_2^0)^{-1}(A'_4A'_1 + A'_2 - (A_4^{0'}A_1^{0'} + A_2^{0'}), A'_4 - A_4^{0'}) (\bar{\tau}^{0'}\tau_c)^{-1}\rho_{\perp} \end{aligned}$$

We replace $A'_2 - A_2^{0'}$ by $-A_4^{0'}(A'_1 - A_1^{0'})$ and find

$$\begin{aligned} C &\sim -(\beta_2^{0'}\beta_2^0)^{-1}((A'_4 - A_4^{0'})A'_1, A'_4 - A_4^{0'}) (\bar{\tau}^{0'}\tau_c)^{-1}\rho_{\perp} \\ &= -(\beta_2^{0'}\beta_2^0)^{-1}(A'_4 - A_4^{0'})(A'_1, I_s) (\bar{\tau}^{0'}\tau_c)^{-1}\rho_{\perp}. \end{aligned}$$

We next want to show that $\rho'_{\perp} = (A'_1, I_s)(\bar{\tau}^{0'}\tau_c)^{-1}$. We find from $\beta_c = \tau_c\rho_c$, that

$$\rho_c = (c, c_1)'\beta_c = \begin{pmatrix} I_r \\ -A'_1 \end{pmatrix}, \quad \rho_{c\perp} = \begin{pmatrix} A_1 \\ I_s \end{pmatrix}.$$

Similarly we find from $\beta = \tau\rho$, by multiplying by $(c, c_1)'$ that

$$(I_r, -A_1)(\bar{\beta}^{0'}\beta_c)^{-1} = (c, c_1)'\beta_c(\bar{\beta}^{0'}\beta_c)^{-1} = (c, c_1)'\tau_c(\bar{\tau}^{0'}\tau_c)^{-1}\rho = (\bar{\tau}^{0'}\tau_c)^{-1}\rho,$$

and hence that

$$\rho_{\perp} = (\tau_c'\bar{\tau}^0)^{-1} \begin{pmatrix} A_1 \\ I_s \end{pmatrix}.$$

Thus we find

$$(A'_1, I_s)(\bar{\tau}^{0'}\tau_c)^{-1}\rho_\perp = \rho'_\perp\rho_\perp \sim \rho_\perp^{0'}\rho_\perp^0$$

and the expression for C becomes

$$C \sim -(\beta_2^{0'}\beta_2^0)^{-1}(A'_4 - A_4^{0'})\rho_\perp^{0'}\rho_\perp^0.$$

Finally we find B_0 : The relation between A_3 and $B_0 = \bar{\beta}_2^{0'}(\psi - \psi^0)$ is given by

$$\begin{aligned} A_3 &= -\bar{\alpha}'\Gamma\beta_2 = \bar{\alpha}'(\alpha\psi' + \bar{\alpha}_\perp\kappa'\tau')\beta_2 = \psi'\beta_2 \\ &= (\psi - \psi^0)'\beta_2^0 + \psi'(\beta_2^0 - \beta_2) + \psi^{0'}\beta_2^0 \\ &= B_0'(\beta_2^{0'}\beta_2^0) + \psi'(\beta_2^0 - \beta_2) + A_3^0, \end{aligned}$$

so that

$$\begin{aligned} B_0 &= (\beta_2^{0'}\beta_2^0)^{-1} [(A'_3 - A_3^{0'}) + (\beta_2 - \beta_2^0)'\psi] \\ &\sim (\beta_2^{0'}\beta_2^0)^{-1} [(A'_3 - A_3^{0'}) - c(A_1A_4 + A_2 - A_1^0A_4^0 - A_2^0) - c_1(A_4 - A_4^0)'\psi^0] \end{aligned}$$

Replacing $A_2 - A_2^0$ by $-(A_1 - A_1^0)A_4^0$ we find

$$\begin{aligned} B_0 &\sim (\beta_2^{0'}\beta_2^0)^{-1} [(A'_3 - A_3^{0'}) - (cA_1(A_4 - A_4^0) + c_1(A_4 - A_4^0))'\psi^0] \\ &\sim (\beta_2^{0'}\beta_2^0)^{-1} [(A'_3 - A_3^{0'}) - (A_4 - A_4^0)'(A_1^{0'}c' + c_1')\psi^0]. \end{aligned}$$

■

2 A simple example.

Let us consider a very stylized example to see what the conditions imply. We let θ_2 bivariate and θ_1 and θ_0 be univariate. We define

$$\begin{aligned} \theta_0(\phi_0, \phi_1, \phi_2) &= \phi_0 \\ \theta_1(\phi_0, \phi_1, \phi_2) &= \phi_1 + d\phi_2 \\ \theta_2(\phi_0, \phi_1, \phi_2) &= \begin{cases} \phi_2 + a\phi_1^2 + b\phi_1 \\ c\phi_0^2 \end{cases} \end{aligned}$$

We assume that $\phi^0 = 0$, and consider the behaviour of θ_i under the local alternative $\phi_0 = T^{-1/2}\eta_0, \phi_1 = T^{-1}\eta_1, \phi_2 = T^{-2}\eta_2$. We normalize θ_i and find

$$\begin{aligned} T^{1/2}\theta_0(T^{-1/2}\eta_0, T^{-1}\eta_1, T^{-2}\eta_2) &= \eta_0 \\ T\theta_1(T^{-1/2}\eta_0, T^{-1}\eta_1, T^{-2}\eta_2) &= \eta_1 + dT^{-1}\eta_2 \\ T^2\theta_2(T^{-1/2}\eta_0, T^{-1}\eta_1, T^{-2}\eta_2) &= \begin{cases} \eta_2 + a\eta_1^2 + Tb\eta_1 \\ Tc\eta_0^2 \end{cases} \end{aligned}$$

Obviously the limit only exists if $b = \partial(\theta_2)_1/\partial\phi_1 = 0$ and $c = \frac{1}{2}\partial^2(\theta_2)_2/\partial\phi_0^2 = 0$. The full set of conditions are collected in (20) and (21). Note that it had no consequences

that θ_1 depended on ϕ_2 , because the terms with $\partial\theta_1/\partial\phi_2$ disappeared in the limit of a local alternative. If in fact $a = b = 0$, then the limit becomes

$$D_T^{-1} \cdot \theta(D_T \cdot \eta) \rightarrow (\eta_0, \eta_1, \eta_2 + a\eta_1^2, 0)'$$

Obviously the limit of the likelihood function, which is quadratic in θ_2 , will be of degree four in η_1 , and if we want a quadratic likelihood function in the limit, we have to assume that also $a = \frac{1}{2}\partial^2(\theta_2)_1/\partial\phi_1^2 = 0$, where the full set of conditions is collected in (21).

Note that the conditions on a and b can always be achieved by a reparametrization. Simply define

$$\tilde{\phi}_2 = \phi_2 + a\phi_1^2 + b\phi_1, \quad \tilde{\phi}_1 = \phi_1, \quad \tilde{\phi}_0 = \phi_0.$$

Thus for this particular example we need $a = b = 0$, in order to achieve a finite limit of the likelihood function under the local alternative and $c = 0$ to get a limit that is quadratic.

In order to investigate the condition (16), we have to evaluate the derivative, not at the true value but at the local value given by the $(T^{-1/2}\psi_0, T^{-1}\psi_1, T^{-2}\psi_2)$. Let us first find the derivative of $\theta = (\theta_0, \theta_1, \theta_{21}, \theta_{22})'$ at the point $\phi = (\phi_0, \phi_1, \phi_2)'$ in the direction $D_T \cdot \eta$:

$$\frac{d\theta}{d\phi'}(D_T^{-1} \cdot \eta) = \begin{pmatrix} \eta_0 \\ \eta_1 + dT^{-1}\eta_2 \\ (2a\phi_1 + b)T\eta_1 + \eta_2 \\ 2c\phi_0 T^{3/2}\eta_0 \end{pmatrix}$$

Hence at the point $(\phi_0, \phi_1, \phi_2) = (T^{-1/2}\psi_0, T^{-1}\psi_1, T^{-2}\psi_2)$ we find

$$D_T^{-1} \cdot \frac{d\theta}{d\phi'}(D_T \cdot \eta) = \begin{pmatrix} \eta_0 \\ \eta_1 + dT^{-1}\eta_2 \\ (2a\psi_1 + Tb)\eta_1 + \eta_2 \\ 2cT\psi_0\eta_0 \end{pmatrix}$$

We see again that in order for the limit to exist, we need the conditions $b = c = 0$, and for the limit to be independent of ψ we must have $a = 0$.

Thus for this simple example we get that condition (16) is equivalent to the conditions given in (20) and (21)

References

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