

## 2. FOURIER EXPANSIONS IN HIGHER DIMENSIONS

**2.1 Multiple Fourier series.**

The theory of Fourier expansions extends readily to higher dimensions. Here the complex formulation is advantageous, because it gives simpler formulas (allowing a better overview than when multiple products of cosines and sines occur everywhere).

Before presenting this, let us underline the fact that is put forward in [A04, Section 2.5], that any  $2p$ -periodic function  $f$  that is *square integrable* on the interval  $[-p, p]$  can be expanded in a Fourier series, with coefficients determined by the Euler formulas on page 39. Moreover, the Bessel inequality and Parseval identity hold for  $f$ . Special cases are piecewise continuous functions, or just bounded (measurable) functions. It can be seen directly from the Euler formulas that the Fourier coefficients are bounded in  $n$ , but the Bessel inequality gives a still better information, namely that  $a_n \rightarrow 0$  and  $b_n \rightarrow 0$  for  $n \rightarrow \infty$ .

Now recall the complex formulation in one variable: It is based on the Euler identity, for  $x \in \mathbb{R}$ :

$$(2.1) \quad \begin{aligned} e^{ix} &= \cos x + i \sin x, & \text{hence} \\ \cos x &= \frac{1}{2}(e^{ix} + e^{-ix}), & \sin x = \frac{1}{2i}(e^{ix} - e^{-ix}). \end{aligned}$$

In the Fourier series of a  $2p$ -periodic function  $f(x)$ ,

$$(2.2) \quad f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos(\frac{n\pi}{p}x) + b_n \sin(\frac{n\pi}{p}x)),$$

we can insert the replacements

$$(2.3) \quad \begin{aligned} \cos(\frac{n\pi}{p}x) &= \frac{1}{2}(e^{i\frac{n\pi}{p}x} + e^{-i\frac{n\pi}{p}x}), \\ \sin(\frac{n\pi}{p}x) &= \frac{1}{2i}(e^{i\frac{n\pi}{p}x} - e^{-i\frac{n\pi}{p}x}); \end{aligned}$$

then

$$(2.4) \quad s_N(x) = a_0 + \sum_{n=1}^N (a_n \cos(\frac{n\pi}{p}x) + b_n \sin(\frac{n\pi}{p}x)) = \sum_{m=-N}^N c_m e^{i\frac{m\pi}{p}x},$$

with

$$(2.5) \quad c_0 = a_0, \quad c_n = \frac{1}{2}(a_n - ib_n), \quad c_{-n} = \frac{1}{2}(a_n + ib_n).$$

This justifies writing (2.2) as

$$(2.6) \quad f(x) = \sum_{n=-\infty}^{\infty} c_n e^{i\frac{n\pi}{p}x}.$$

One has that

$$(2.7) \quad c_n = \frac{1}{2p} \int_{-p}^p f(x) e^{-i\frac{n\pi}{p}x} dx \text{ for all } n \in \mathbb{Z},$$

which holds also when complex-valued functions  $f(x)$  are allowed (still with  $x \in \mathbb{R}$ ). The Parseval identity is:

$$(2.8) \quad \sum_{n=-\infty}^{\infty} |c_n|^2 = \frac{1}{2p} \int_{-p}^p |f(x)|^2 dx.$$

Theorem 1.2 says in the complex formulation that when  $f$  is PC1C with period  $2p$ , then:

$$(2.9) \quad \begin{aligned} \text{(i)} \quad & c_n(f') = i\frac{n\pi}{p}c_n(f) \text{ for all } n \in \mathbb{Z}, \\ \text{(ii)} \quad & \sum_{n=-\infty}^{\infty} |c_n| < \infty, \end{aligned}$$

and (iii) the Fourier series converges uniformly (and absolutely) to  $f$ .

It is not hard to extend the ideas to higher dimensions. For simplicity in the formulas we now let  $p = \pi$  and leave to the reader to do the scaling when other lengths are needed.

On  $\mathbb{R}^k$  with points denoted  $x = (x_1, \dots, x_k)$  we consider functions  $f(x)$  that have period  $2\pi$  in each variable  $x_1, \dots, x_k$ . They are completely determined by their values on the cube  $[-\pi, \pi]^k$ . The elements of  $\mathbb{Z}^k$  will be denoted  $\mathbf{n} = (n_1, \dots, n_k)$ , with length

$$(2.10) \quad \|\mathbf{n}\| = \sqrt{n_1^2 + \dots + n_k^2}.$$

The functions

$$(2.11) \quad e^{i\mathbf{n}\cdot x} = e^{i(n_1x_1 + \dots + n_kx_k)}, \quad \mathbf{n} \in \mathbb{Z}^k,$$

are  $2\pi$ -periodic in each variable  $x_j$  and satisfy

$$(2.12) \quad (e^{i\mathbf{n}\cdot x}, e^{i\mathbf{m}\cdot x}) = \begin{cases} 0 & \text{if } \mathbf{n} \neq \mathbf{m}, \\ (2\pi)^k & \text{if } \mathbf{n} = \mathbf{m}, \end{cases}$$

when we use the scalar product (inner product)

$$\begin{aligned} (f, g) &= \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \dots \int_{-\pi}^{\pi} f(x_1, x_2, \dots, x_k) \bar{g}(x_1, x_2, \dots, x_k) dx_1 dx_2 \dots dx_k \\ &= \int_{[-\pi, \pi]^k} f(x) \bar{g}(x) dx. \end{aligned}$$

For,

$$(e^{i\mathbf{n}\cdot x}, e^{i\mathbf{m}\cdot x}) = \int_{-\pi}^{\pi} e^{in_1x_1} e^{-im_1x_1} dx_1 \cdots \int_{-\pi}^{\pi} e^{in_kx_k} e^{-im_kx_k} dx_k;$$

here if  $n_j \neq m_j$  for some  $j$ , the integral in  $x_j$  gives a factor 0; on the other hand if  $n_j = m_j$  for all  $j$ , each integral over  $[-\pi, \pi]$  contributes with a factor  $2\pi$ .

It is shown on the basis of the one-dimensional result that a square integrable function  $f(x)$  on  $[-\pi, \pi]^k$ , extended to be  $2\pi$ -periodic in each variable  $x_1, \dots, x_k$ , can be expanded in a Fourier series

$$(2.13) \quad \begin{aligned} f(x) &\sim \sum_{\mathbf{n} \in \mathbb{Z}^k} c_{\mathbf{n}} e^{i\mathbf{n}\cdot x}, \text{ where} \\ c_{\mathbf{n}} &= \frac{1}{(2\pi)^k} \int_{[-\pi, \pi]^k} f(x) e^{-i\mathbf{n}\cdot x} dx, \end{aligned}$$

in such a way that the partial sum

$$(2.14) \quad s_N(x) = \sum_{\max\{|n_1|, \dots, |n_k|\} \leq N} c_{\mathbf{n}} e^{i\mathbf{n}\cdot x}$$

converges in the mean to  $f(x)$ , in the sense that

$$(2.15) \quad \int_{[-\pi, \pi]^k} |f(x) - s_N(x)|^2 dx \rightarrow 0 \text{ for } N \rightarrow \infty.$$

Here the following Parseval identity holds:

$$(2.16) \quad \sum_{\mathbf{n} \in \mathbb{Z}^k} |c_{\mathbf{n}}|^2 = \frac{1}{(2\pi)^k} \int_{[-\pi, \pi]^k} |f(x)|^2 dx.$$

As a corollary to the Parseval identity we see that  $|c_{\mathbf{n}}| \rightarrow 0$  for  $\|\mathbf{n}\| \rightarrow \infty$ ; this holds under the mere assumption that  $f$  is square integrable on  $[-\pi, \pi]^k$ . We give below some information on uniform convergence.

For  $k = 2$ , the formulation with cosine and sine is found from the above by noting that

$$e^{i(n_1x_1+n_2x_2)} = (\cos n_1x_1 + i \sin n_1x_1)(\cos n_2x_2 + i \sin n_2x_2).$$

For  $n_1$  and  $n_2 \in \mathbb{N}$  we can use this in the four terms

$$\begin{aligned} c_{(n_1, n_2)} e^{i(n_1x_1+n_2x_2)} + c_{(n_1, -n_2)} e^{i(n_1x_1-n_2x_2)} \\ + c_{(-n_1, n_2)} e^{i(-n_1x_1+n_2x_2)} + c_{(-n_1, -n_2)} e^{i(-n_1x_1-n_2x_2)} \end{aligned}$$

and regroup them as a linear combination of  $\cos n_1x_1 \cos n_2x_2$ ,  $\cos n_1x_1 \sin n_2x_2$ ,  $\sin n_1x_1 \cos n_2x_2$  and  $\sin n_1x_1 \sin n_2x_2$ . This is somewhat unmanageable, but it becomes more manageable when we restrict the attention to functions that are *odd* in  $x_1$  as well as  $x_2$ ; they only have sine terms

$$(2.17) \quad \begin{aligned} f(x) &\sim \sum_{n_1, n_2 \in \mathbb{N}} b_{n_1, n_2} \sin n_1x_1 \sin n_2x_2, \text{ with} \\ b_{n_1, n_2} &= -c_{(n_1, n_2)} + c_{(n_1, -n_2)} + c_{(-n_1, n_2)} - c_{(-n_1, -n_2)}, \end{aligned}$$

since, in the calculation of (2.13), a cosine factor  $\cos n_i x_i$  integrated together with  $f$  in the  $x_i$ -variable gives 0. Here

$$b_{n_1, n_2} = \frac{4}{\pi^2} \int_0^\pi \int_0^\pi f(x_1, x_2) \sin n_1 x_1 \sin n_2 x_2 dx_1 dx_2.$$

Note however that when one differentiates a sine series, cosine comes in again.

There is a general result:

**Theorem 2.1.**

1° If  $f(x)$  is  $2\pi$ -periodic in each coordinate, and  $C^1$ , then for all  $\mathbf{n} \in \mathbb{Z}^k$ ,

$$(2.18) \quad c_{\mathbf{n}}\left(\frac{\partial f}{\partial x_j}\right) = i n_j c_{\mathbf{n}}(f), \quad j = 1, \dots, k,$$

and the Parseval identity for the derivatives implies

$$(2.19) \quad \sum_{\mathbf{n} \in \mathbb{Z}^k} \|\mathbf{n}\|^2 |c_{\mathbf{n}}(f)|^2 < \infty.$$

Moreover, if  $f$  is  $C^l$  for some  $l \geq 1$ , then

$$(2.20) \quad \sum_{\mathbf{n} \in \mathbb{Z}^k} \|\mathbf{n}\|^{2l} |c_{\mathbf{n}}(f)|^2 < \infty.$$

2° For  $k = 2$  or  $3$ , if  $f(x)$  is  $2\pi$ -periodic in each coordinate and  $C^{l+2}$ , then

$$(2.21) \quad \sum_{\mathbf{n} \in \mathbb{Z}^k} \|\mathbf{n}\|^l |c_{\mathbf{n}}(f)| < \infty.$$

The estimate (2.21) implies that the Fourier series and its termwise differentiated series up to order  $l$  are uniformly convergent.

*Indications of proof.* In 1°, the identity in (2.18) is shown by integration by parts (in the  $x_j$ -variable) in the formula for  $c_{\mathbf{n}}\left(\frac{\partial f}{\partial x_j}\right)$ . Then the Parseval identity for  $\frac{\partial f}{\partial x_j}$  implies the convergence of the series  $\sum_{\mathbf{n}} |n_j|^2 |c_{\mathbf{n}}(f)|^2$ . When we sum over  $j$  we find (2.19). When  $f$  is  $C^l$ , this can be applied for any succession of  $l$  partial derivatives, showing that

$$\sum_{\mathbf{n} \in \mathbb{Z}^k} |p(n_1, \dots, n_k)|^2 |c_{\mathbf{n}}(f)|^2 < \infty$$

for any polynomial  $p$  of degree  $l$ . Since  $\|\mathbf{n}\|^{2l}$  is bounded by a sum of squares of such polynomials, the result (2.20) follows.

For 2°, note that it is here a question of series with  $|c_{\mathbf{n}}|$  in the first power only. One can show that the series  $\sum_{\mathbf{n} \in \mathbb{Z}^k \setminus \{\mathbf{0}\}} \|\mathbf{n}\|^{-4}$  is convergent for  $k = 2$  and  $3$  (this is related to

the fact that  $\int_{|x| \geq 1} |x|^{-4} dx < \infty$  in dimensions  $< 4$ ; see also Exercise E2.3). Then we can do a trick as in the proof of Theorem 1.2:

$$\begin{aligned} \sum_{\|\mathbf{n}\| \leq N} \|\mathbf{n}\|^l |c_{\mathbf{n}}| &= \sum_{0 < \|\mathbf{n}\| \leq N} \|\mathbf{n}\|^{-2} \|\mathbf{n}\|^{l+2} |c_{\mathbf{n}}| \\ &\leq \frac{1}{2} \left( \sum_{0 < \|\mathbf{n}\| \leq N} \|\mathbf{n}\|^{-4} + \sum_{0 < \|\mathbf{n}\| \leq N} \|\mathbf{n}\|^{2(l+2)} |c_{\mathbf{n}}|^2 \right) \leq C, \end{aligned}$$

for all  $N$ , where we for the last series use that (2.20) holds with  $l$  replaced by  $l + 2$ .

The estimate (2.21) implies uniform convergence of the termwise differentiated series up to order  $l$ , since  $|e^{i\mathbf{n} \cdot x}| = 1$ , and

$$(2.22) \quad \frac{\partial}{\partial x_j} e^{i\mathbf{n} \cdot x} = i n_j e^{i\mathbf{n} \cdot x},$$

so (2.21) (times a constant) is a majorizing series for all those termwise derived series.  $\square$

**Remark 2.2.**

(a) There is a result along the lines of 2° also in higher dimensions. One can show that for general  $k$ , the series  $\sum_{\mathbf{n} \in \mathbb{Z}^k \setminus \{\mathbf{0}\}} \|\mathbf{n}\|^{-2k'}$  is convergent when  $2k' > k$ . Then when the Fourier coefficients of  $f$  satisfy the estimate (2.20) with  $l$  replaced by  $l + k'$ , they will satisfy (2.21), by a version of the above trick.

This is actually an example of Sobolev's Theorem — prominent in the more advanced theory — that says that functions with finite Sobolev norm of order  $l + k'$ , some  $k' > k/2$ , are in  $C^l$  (here the squareroot of (2.20) plays the role of the  $l$ 'th Sobolev norm).

(b) In part 1° of Theorem 2.1, the formulas (2.18)–(2.19) can be shown under slightly weaker assumptions. It suffices that  $f$  is continuous with square integrable first derivatives defined in some reasonable sense. For example, if  $[-\pi, \pi]^k$  is divided into a finite number of polyedric subdomains, and the first derivatives of  $f$  are defined in each of these subdomains and extend to continuous functions on their closures, then (2.18) and (2.19) hold. We can call such derivatives *piecewise continuous* (although the notion could be defined also in more general situations). Similarly, if  $f$  is  $C^{l-1}$  and the  $l$ 'th order derivatives are piecewise continuous, then (2.20) holds.

In Theorem 2.1 2°, it is then sufficient for (2.21) that  $f$  is  $C^{l+1}$  with piecewise continuous derivatives of order  $l + 2$ .

**2.2 The wave equation with initial data on a rectangle.**

Using Theorem 2.1, we can justify the solution formula for the two-dimensional wave equation in [A04, Section 3.7] as follows (using the notation  $(x, y)$  for a point in  $\mathbb{R}^2$ ):

**Theorem 2.3.** 1° *When  $f(x, y)$  is  $C^2$  and  $g(x, y)$  is  $C^1$  on  $M = [0, a] \times [0, b]$ , and  $f$  and  $g$  are zero on the boundary  $\partial M = \{(x, y) \in M \mid x = 0 \text{ or } a, y = 0 \text{ or } b\}$ , then*

$$(2.23) \quad \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} (|B_{mn}| + |B_{mn}^*|) < \infty.$$

*Then the series in (4) converges uniformly on  $M \times [0, \infty[$  to a continuous function  $u(x, y, t)$  satisfying the boundary condition  $u = 0$  for  $(x, y) \in \partial M$  and the first initial condition  $u = f$  for  $t = 0$ .*

2° When furthermore  $f(x, y)$  is  $C^4$  and  $g(x, y)$  is  $C^3$  on  $M$ , and

$$(2.24) \quad \frac{\partial^2 f}{\partial x^2}, \frac{\partial^2 f}{\partial y^2}, \frac{\partial^2 g}{\partial x^2} \text{ and } \frac{\partial^2 g}{\partial y^2} \text{ are 0 on } \partial M,$$

then

$$(2.25) \quad \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} (m^2 + n^2)(|B_{mn}| + |B_{mn}^*|) < \infty,$$

and  $u(x, y, t)$  is  $C^2$  on  $M \times [0, \infty[$  and satisfies the wave equation and the initial- and boundary conditions.

*Proof.* When  $f$  and  $g$  satisfy the hypotheses in 1°, they extend to functions  $f^*$  resp.  $g^*$  on  $\mathbb{R}^2$  that are odd in  $x$  with period  $2a$ , and odd in  $y$  with period  $2b$ , such that the extended functions are in  $C^1$  on  $\mathbb{R}^2$  and the second derivatives of  $f^*$  are piecewise continuous. The  $B_{mn}$  are the coefficients in the sine expansion of  $f$  in two variables as in (2.17), so by Theorem 2.1 2° with  $l = 0$  and Remark 2.2(b), the series  $\sum_{m,n} |B_{mn}|$  is convergent. The  $B_{mn}^*$  satisfy

$$(2.26) \quad B_{mn}^* = \frac{b_{mn}^*}{\lambda_{mn}},$$

where the  $b_{mn}^*$  are the coefficients in the sine expansion of  $g$  in two variables. By 1° of Theorem 2.1,

$$(2.27) \quad \sum_{m,n \in \mathbb{N}} (m^2 + n^2) |b_{mn}^*|^2 < \infty.$$

Since  $\lambda_{mn} = c\pi \sqrt{\frac{m^2}{a^2} + \frac{n^2}{b^2}}$  clearly satisfies

$$(2.28) \quad c_1(m^2 + n^2)^{\frac{1}{2}} \leq \lambda_{mn} \leq c_2(m^2 + n^2)^{\frac{1}{2}}$$

with positive constants  $c_1$  and  $c_2$ , we conclude from (2.26) and (2.27) that the series of  $|B_{mn}^*|^2$  satisfies

$$(2.29) \quad \sum_{m,n \in \mathbb{N}} (m^2 + n^2)^2 |B_{mn}^*|^2 < \infty.$$

From this we deduce as in the proof of Theorem 2.1 2° that

$$(2.30) \quad \sum_{m,n \in \mathbb{N}} |B_{mn}^*| < \infty.$$

This completes the proof of (2.23), which implies uniform convergence as stated.

When furthermore the hypotheses of 2° are satisfied, the extensions  $f^*$  and  $g^*$  are  $C^3$  on  $\mathbb{R}^2$  and the fourth-order derivatives of  $f^*$  are piecewise continuous. It follows from Theorem 2.1 2° and Remark 2.2(b) that

$$(2.30) \quad \sum_{m,n} (m^2 + n^2) |B_{mn}| < \infty, \quad \sum_{m,n} (m^2 + n^2)^{\frac{1}{2}} |b_{mn}^*| < \infty.$$

Using again (2.26) and (2.28) we conclude that (2.25) holds. Then the termwise differentiated series up to order 2 converge uniformly on  $M \times [0, \infty[$ , since each differentiation in  $x$  or in  $y$  essentially gives a factor  $n$  or  $m$ , and each differentiation in  $t$  gives a factor  $\lambda_{mn}$ . Thus the differential equation and the remaining initial condition can be verified.  $\square$

The conditions on  $f$  and  $g$  in the theorem are *sufficient conditions* — one can weaken them a little and still get solvability — but at least they give some firm ground for the claim that the described procedure gives a solution of the problem posed. (For example, in Theorem 2.3 2°, one can allow the fourth-order derivatives of  $f$  and the third-order derivatives of  $g$  to be piecewise continuous on  $M$ .)

In Example 3.7.1,  $h(x) = x(1 - x)$  extends to an odd, 2-periodic function  $h^*$  whose first derivative is a continuous triangular function, and the second derivative is piecewise constant, having jumps at the period points  $2n$ ,  $n \in \mathbb{Z}$ . Then the odd, 2-periodic extension of  $f(x, y)$  is  $C^1$  with piecewise continuous (in fact piecewise constant) second derivatives. Here Theorem 2.3 1° gives that the series for  $u(x, y, t)$  converges uniformly (which is also clear from the formulas), but (2.24) is not satisfied, and the differential equation holds only in a generalized sense. (In a deeper analysis one can show that the solution is smooth in large areas, but that the irregularities in the initial value propagate along characteristic cones when  $t$  increases.)

### 2.3 The heat equation with initial data on a rectangle.

The solution formulas for the heat problem [A04, page 161] can be checked in a similar way. Here we find, as for the one-dimensional heat equation, that the solution becomes  $C^\infty$  as soon as  $t$  becomes positive.

**Theorem 2.4.** 1° *When  $f(x, y)$  is  $C^2$  on  $M = [0, a] \times [0, b]$ , and is zero on the boundary  $\partial M$ , then*

$$(2.31) \quad \sum_{m,n \in \mathbb{N}} |A_{mn}| < \infty,$$

*and the series in (13) converges uniformly on  $M \times [0, \infty[$  to a continuous function  $u(x, y, t)$  satisfying the boundary condition  $u = 0$  for  $(x, y) \in \partial M$  and the initial condition  $u = f$  for  $t = 0$ .*

2° *Assume merely that  $f(x, y)$  is square integrable on  $M$ . Then the series in (13) and all the termwise differentiated series of arbitrarily high order converge uniformly on  $M \times [\varepsilon, \infty[$ , for any  $\varepsilon > 0$ . In particular, the differential equation (11) and the boundary condition (12) are verified for  $t > 0$ .*

*Proof.* Part 1° is shown in the same way as in Theorem 2.3; the expansion coefficients of  $f$  are now called  $A_{mn}$ , and the hypotheses assure that the odd periodic extension of  $f$  is  $C^1$

with piecewise continuous second derivatives, so that (2.31) holds and defines a majorizing series for (13).

For part 2°, we observe that when  $t \geq \varepsilon$ , then

$$(2.32) \quad e^{-\lambda_{mn}^2 t} \leq e^{-\lambda_{mn}^2 \varepsilon} \leq e^{-c_1^2 \varepsilon (m^2 + n^2)},$$

cf. (2.28). For any  $j, k, l \geq 0$ , application of  $\frac{\partial^j}{\partial x^j} \frac{\partial^k}{\partial y^k} \frac{\partial^l}{\partial t^l}$  termwise gives a series

$$(2.33) \quad \pm \sum_{m, n \in \mathbb{N}} A_{mn} \left(\frac{m\pi}{a}\right)^j \left(\frac{n\pi}{b}\right)^k \lambda_{mn}^{2l} \frac{\sin}{\cos} \left(\frac{m\pi}{a} x\right) \frac{\sin}{\cos} \left(\frac{n\pi}{b} y\right) e^{-\lambda_{mn}^2 t}.$$

Since the  $|A_{mn}|$  are bounded by a constant (cf. (2.16)ff.), this is majorized by a convergent series

$$(2.34) \quad \sum_{m, n \in \mathbb{N}} c'(m^2 + n^2)^{\frac{1}{2}j + \frac{1}{2}k + l} e^{-c_1^2 \varepsilon (m^2 + n^2)} < \infty,$$

where the convergence follows e.g. since  $(m^2 + n^2)^{\frac{1}{2}j + \frac{1}{2}k + l} e^{-c_1^2 \varepsilon (m^2 + n^2)} \leq c''(m^2 + n^2)^{-2}$ . Thus  $u$  is  $C^\infty$  for  $t > 0$ .  $\square$

When  $f$  is merely square integrable, one can say that the initial condition is verified in the sense that the series (13) for  $t = 0$  converges in the mean to  $f$  on  $M$ .

## 2.4 The Poisson equation with zero boundary data.

For the solution of the Poisson equation on a rectangle put forward in [A04, Section 3.9], sufficient conditions for convergence can likewise be found from Theorem 2.1 ff. The discussion goes rather similarly to that in Theorems 2.3–4:

When the odd periodic extension  $f^*$  of  $f$  is  $C^1$  with piecewise continuous second derivatives, the sinus coefficients  $A_{mn}$  of  $f$  satisfy (2.31). Then since  $E_{mn} = -A_{mn}/\lambda_{mn}$ , one has that

$$(2.35) \quad \sum_{m, n \in \mathbb{N}} (m^2 + n^2) |E_{mn}| < \infty,$$

so the differential equation can be verified by termwise differentiation, and  $u(x, y)$  is indeed a solution of the problem; it lies in  $C^2(M)$ .

The solution found in Example 3.9.1 is a generalized solution, and the calculations in Example 3.9.2 are quite formal (the series for  $u$  does not satisfy our criteria for termwise differentiation of order 2).

In more advanced treatments of the various differential equations, the theory of Sobolev spaces provides a more satisfactory framework than the spaces of  $C^l$ -functions.

## 2.5 Convergence analysis for the Laplace equation on a disk.

Consider the series solution established in [A04, Section 4.4] for the Laplace equation on a disk with radius  $a$ , with a prescribed boundary value  $f$ . We shall show that when  $f$  is merely square integrable, the series and all termwise differentiated series converge uniformly on the disks with the same center and radius  $< a$ . For simplicity in the formulation, let us take  $a = 1$  and leave the scaling to the general case to the reader.



When  $f(\theta)$  is square integrable on  $[0, 2\pi]$ , its Fourier coefficients are bounded:

$$(2.36) \quad |a_n| \leq C, \quad |b_n| \leq C, \quad \text{for all } n.$$

Let  $\varepsilon \in ]0, 1[$ . Then the series for  $u$ ,

$$(2.37) \quad u(r, \theta) = a_0 + \sum_{n \in \mathbb{N}} r^n (a_n \cos n\theta + b_n \sin n\theta)$$

has the majorizing series, when  $r \leq 1 - \varepsilon$ :

$$(2.38) \quad \sum_{n \geq 0} C(1 - \varepsilon)^n.$$

We can check termwise derivatives in  $r$  and  $\theta$ , showing that each differentiation essentially gives a factor  $n$ , so that the series after  $k$  differentiations is majorized by a series

$$(2.39) \quad \sum_{n \geq 0} C'(1 + n)^k (1 - \varepsilon)^n.$$

Why is this convergent? Apply for example the quotient criterion, or note that  $(1 - \varepsilon)^n = e^{-sn}$ , where  $-s = \ln(1 - \varepsilon) < 0$ ; here since the exponential function wins over any polynomial,  $(1 + n)^k e^{-sn} \leq C''(1 + n)^{-2}$  (as we have used before).

However, if we treat the  $(r, \theta)$ -derivatives of  $u$ , we still have to worry about how the information carries over to the  $(x, y)$ -derivatives (in the original coordinates), especially how things fit together at  $r = 0$ . But there is another point of view that gives the  $(x, y)$ -behavior directly:

When the Fourier series of  $f$  is written in the complex form

$$(2.40) \quad f(\theta) = \sum_{m=-\infty}^{\infty} c_m e^{im\theta},$$

we get the series for  $u$  in the complex form, that we can reformulate further:

$$(2.41) \quad \begin{aligned} u(r, \theta) &= \sum_{m \in \mathbb{Z}} r^{|m|} c_m e^{im\theta} \\ &= \sum_{m \geq 0} c_m (r e^{i\theta})^m + \sum_{m' > 0} c_{-m'} (r e^{-i\theta})^{m'} \\ &= \sum_{m \geq 0} c_m (x + iy)^m + \sum_{m' > 0} c_{-m'} (x - iy)^{m'}, \end{aligned}$$

with bounded coefficients. Here, when  $(x, y)$  lies in the disk with radius  $1 - \varepsilon$ ,  $x + iy$  and  $x - iy$  both have absolute value  $\leq 1 - \varepsilon$ . Termwise differentiations in  $x$  and  $y$  give polynomials in  $m$  resp.  $m'$  as factors. Then we can again use series of the type in (2.39) as majorizing series, and find that all termwise derived series are uniformly convergent on the smaller disk. Thus  $u$  is  $C^\infty$  there and satisfies  $\Delta u = 0$ .

## E2. EXERCISES

**Exercise E2.1.** Consider the cosine-sine formulation of the Fourier expansion of a function  $f(x_1, x_2)$  that is  $2\pi$ -periodic in each variable. You are asked to express the coefficients of the functions  $\cos n_1 x_1 \cos n_2 x_2$  and  $\cos n_1 x_1 \sin n_2 x_2$  in terms of the coefficients  $c_{\mathbf{n}}$  in the series in (2.13). What is the constant term?

**Exercise E2.2.** Answer Exercise 3.8.12 in [A04], with the additional point:

(d) Show that the differential equation is verified for  $z < c$ , when  $f(x, y)$  is square integrable.

**Exercise E2.3.** In the following, we identify  $\mathbb{Z}^2$  with a subset of  $\mathbb{R}^2$ , namely with the points with integer coordinates  $(n_1, n_2)$ .

(a) For each  $l \in \mathbb{N}$ , show that there are  $8l$  of these integer points  $(n_1, n_2)$  on the boundary of the square  $[-l, l] \times [-l, l]$ , and that they satisfy

$$n_1^2 + n_2^2 \geq l^2.$$

(b) Show that there is a constant  $c$  such that

$$\sum_{|n_1| \leq l, |n_2| \leq l, (n_1, n_2) \neq (0, 0)} \frac{1}{(n_1^2 + n_2^2)^2} \leq \sum_{1 \leq j \leq l} \frac{c}{j^3}$$

(c) Show that the series

$$\sum_{(n_1, n_2) \in \mathbb{Z}^2 \setminus \{(0, 0)\}} \frac{1}{(n_1^2 + n_2^2)^2}$$

is convergent.